

# Model for WDM Subhalo Distribution

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### **Power Spectrum of WDM**



The matter power spectrum of WDM will be suppressed at small scale compared with CDM case.



# **Simulation and Spurious Halos**



10<sup>13</sup>

10<sup>16</sup>

 $10^{15}$ 

1014

 $M_{h}(h^{-1}M_{\odot})$ 

1.0 0.5 0.0

10<sup>12</sup>





#### Warm Dark Matter



# **Model Framework**

The spatial distribution of the subhalo should trace the halo density profile.

"anti-bias": subhalo final number density is less centrally concentrated





# **Model Framework**



# **Unevolved spatial distribution**

The number density profile of accreted subhalos traces the halo density profile.

$$ilde{n}_{
m sub}(R|m_{
m acc}) \sim { ilde{
ho}}_{
m DM}(R)$$
 .

Unevolved spatial distribution

Halo density profile









The subhalo mass function of WDM is suppressed at the scale below  $M_{\rm hm}$  in

$$rac{n_X(M)}{n_{ ext{CDM}}(M)}\simeq \left(1+\left(arac{M_{ ext{hm}}}{M}
ight)^b
ight)^c \quad ext{Stucker et al. 2021}$$

unevolved mass function 
$$a=2.3$$
 evolved mass function  $b=1$ 

c = -0.68



# **Tidal Stripping**



The ratio between the final mass and infall mass of a given subhalo is approximately proportional to the halocentric radius, and the scatter follows the lognormal distribution Survival fraction depends on subhalo mass, more WDM subhaloes below  $M_{\rm hm}$  are disrupted.



# Mass Stripping In WDM



Stronger tidal stripping on WDM subhaloes

Vulnerable to the tidal effects



## **Halo Concentration**



Later formation time of WDM halos leads to lower concentration.

Bose et al. 2016



$$rac{c_{200}^{
m WDM}}{c_{200}^{
m CDM}} = \left(1+\gamma_1rac{M_{
m hm}}{M_{200}}
ight)^{-\gamma_2} imes(1+z)^{eta(z)}$$





# **MCMC Realization**

$m_{\chi}/{ m keV}$	$M_{ m hm}/h^{-1}M_{\odot}$	$\mu_*$	eta	$\sigma$
CDM	~	0.48	1.12	0.99
3.0	$2.3 \times 10^{8}$	0.46	1.2	1.01
1.2	$5.4 \times 10^{9}$	0.42	1.3	1.1
0.5	$1.1 \times 10^{11}$	0.29	1.45	1.25
	data,CDM data,3.0keV data,1.2keV data,0.5keV – model,CDM – model,3.0keV	dN(m, R)	$f^{m_{max}}$	$[m_{m}]^{-\alpha}$



$$egin{aligned} rac{\mathrm{d}N(m,R)}{\mathrm{d}\ln m \ \mathrm{d}^3 R} \sim & ilde{
ho}(R) \int_{m_{\mathrm{min}}}^{m_{\mathrm{max}}} f_s(m_{\mathrm{acc}}) iggl[rac{m_{\mathrm{acc}}}{m_0}iggr]^{-lpha} iggl(1 + iggl(\kappa rac{M_{\mathrm{hm}}}{m_{\mathrm{acc}}}iggr)^\etaiggr)^\gamma \ & imes \expiggl[-rac{1}{2}iggl(rac{\ln \mu - \ln ar{\mu}(R)}{\sigma}iggr)^2iggr] \mathrm{d}\ln m_{\mathrm{acc}} \end{aligned}$$



# **Model Predictions**





#### Summary



- In contrast to the power-law form in CDM, the unevolved subhalo mass function for WDM is suppressed at the low mass end due to the cut-off in the power spectrum.
- WDM subhaloes are more vulnerable to tidal stripping and disruption due to their lower concentrations at accretion time.
- These differences result in a mass-dependent spatial distribution of WDM subhaloes which also depends on the WDM particle mass.